

The WAVE EQUATION -

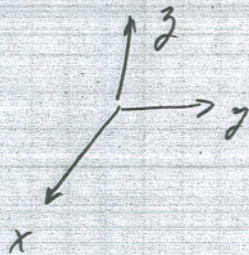
LANGUAGE - ACOUSTICS IMPLIES INVISCID FLUID -

No SHEAR -

USE SLINKY & DIAGRAM

A SOLID OR VERY VISCOUS FLUID HAS $G \neq 0$ -

INTERESTING NOTE AS $E \rightarrow \infty$
 $G \rightarrow 0$



$$\nabla^2 \phi = \frac{1}{c^2} \phi$$

$$\nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$$

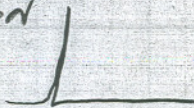
PROPAGATION - UNBOUNDED BODIES

REFLECTION }
REFRACTION } BOUNDARY OR INTERFACE

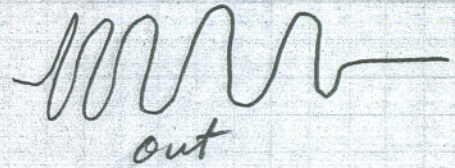
GUIDED WAVES - LIKE 2 BOUNDARIES

1 UNBOUNDED DIMENSION

DISPERSION
GROUP VELOCITY

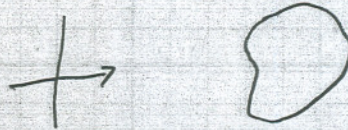


IN



SCATTERING PROBLEMS

→ FINITE DIMENSIONAL BODY



LOOK AT IDEA OF DIMENSIONS!

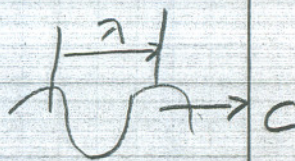
50 SHEETS
22-141
100 SHEETS
22-142
200 SHEETS
22-144



CHARACTERISTIC DIMENSION INTRINSIC

$$\lambda = c/f$$

HOW BIG IS A CAR



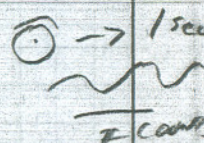
JAPAN MICRO CAR 3 METERS

SPEED OF SOUND IN WATER

1482 m/s

" " " " AIR

332 m/s



WAVE REFERRED TO CHARACTERISTIC LENGTH

5 kHz - $\lambda = .296 M$ WATER
5000 1/s - $\lambda = .066 M$ AIR

MICROCAR IS BIG $> 5\lambda$

Linearized Theory -

$$L(\alpha x + \beta y) = \alpha Lx + \beta Ly$$

Pulses

FOURIER ANALYSIS -

SUPERPOSITION OF HARMONIC WAVES
(TIME HARMONIC WAVES)

HARMONIC $\sin \frac{\pi n t}{T}$ $\cos \frac{\pi n t}{T}$

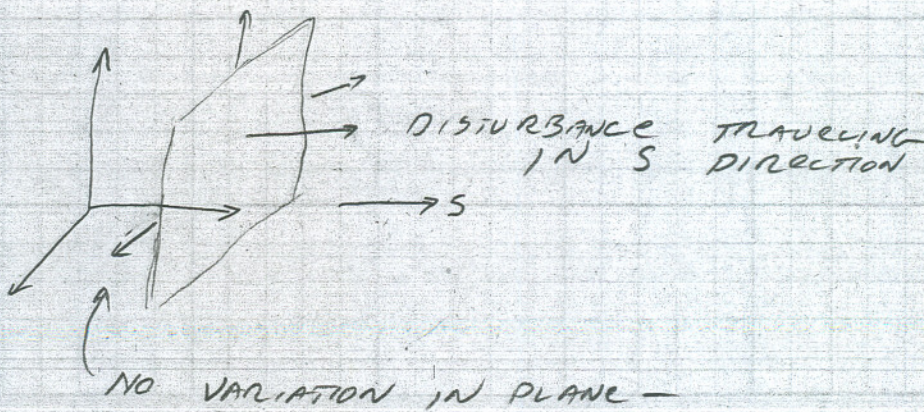
OR Euler's Identity

$$e^{i\theta} = \cos \theta + i \sin \theta$$

START IN 1-DIMENSION -

WAVE ON A STRING -

OR WAVE IN AN UNBOUNDED MEDIA - PLANE WAVE -



50 SHEETS
22-142 100 SHEETS
22-144 200 SHEETS

22-141
22-142
22-144



PARTIAL DIFFERENTIAL EQUATION -

$$\frac{\partial^2 p}{\partial s^2} - \frac{1}{c^2} \frac{\partial^2 p}{\partial t^2} = 0$$

OR $\nabla^2 p + \frac{1}{c^2} \frac{\partial^2 p}{\partial t^2} = 0$

$$\left(\frac{\partial}{\partial s} - \frac{1}{c} \frac{\partial}{\partial t}\right) \left(\frac{\partial}{\partial s} + \frac{1}{c} \frac{\partial}{\partial t}\right) p = 0$$

$p =$ pressure -
VARIES

$$p = p_0 + p'$$

↙ DISTURBANCE
OF PRESSURE
WITH SPATIAL
VARIATION

↑ pressure in FLUID

$$\rho = \rho_0 + \rho'$$
 DENSITY IS PERTURBED
BY PRESSURE VARIATION

MOST GENERAL SOLUTION IS A
PAIR OF DISTURBANCES MOVING
IN THE $+s$ & $-s$ DIRECTION
OF ARBITRARY SHAPE

d'Alembert
Solution

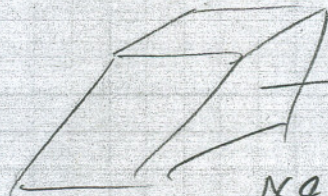
$$p = f\left(t - \frac{1}{c}s\right) + g\left(t + \frac{1}{c}s\right)$$

↘ MOVES IN $+s$

↗ MOVES IN $-s$

VELOCITY OF DISTURBANCES IS c
NO CHANGE IN SHAPE -

TALK ABOUT
LINEAR
THEORY -



→ NET ENERGY
FLUX -

NO CHANGE IN MASS -

WE CAN ALWAYS DESCRIBE AS A
VELOCITY INSTEAD OF A PRESSURE

$$\rho_0 \frac{\partial v}{\partial t} = \frac{\partial p}{\partial t}$$

$$v = \frac{1}{\rho c} \left[f\left(t - \frac{1}{c}s\right) - g\left(t + \frac{1}{c}s\right) \right]$$

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SINGLE FREQUENCY

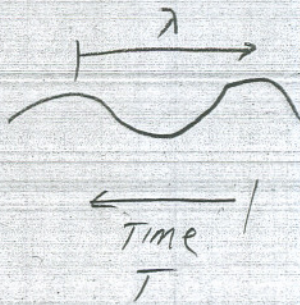
$$\begin{aligned}
 p &= p_{max} \cos(\omega t - \phi) \\
 &= p_{max} \sin(\omega t - \phi') \\
 &= \text{Re } \hat{p} e^{-i\omega t}
 \end{aligned}$$

Not any different
 $f(t) = \sin(\omega t)$
 $= \sin(kx - \omega t)$
 $= \sin(kx - \omega t)$

Same Narration

Remember

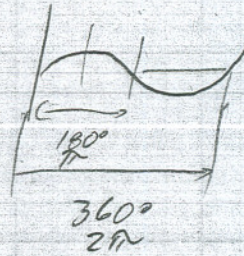
$$T = 1/f$$



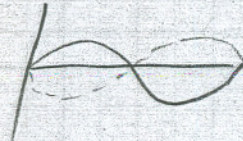
$$\omega = kc$$

k = WAVE NUMBER
 c = VELOCITY

ϕ - PHASE



2 WAVES 180° OUT OF PHASE $\phi = \pi$ RADIANS



Now THE BIG DOG -

SINGLE POINT SOURCE IN AN UNBOUNDED DOMAIN

$$\begin{aligned}
 r^2 &= x_1^2 + x_2^2 + x_3^2 \\
 2r \frac{\partial r}{\partial x_i} &= 2x_i
 \end{aligned}$$

$$\frac{\partial}{\partial x_i} p(r, t) = \frac{\partial p}{\partial r} \frac{\partial r}{\partial x_i} = \frac{x_i}{r} \frac{\partial p}{\partial r}$$

SPHERICAL
WAVE EQUATION

! SPREAD OUTWARD FROM A POINT NO ϕ dependence

$$\frac{1}{r} \frac{\partial^2}{\partial r^2} r p - \frac{1}{c^2} \frac{\partial^2 p}{\partial t^2} = 0$$

PLANE WAVE EXCEPT FOR $\frac{1}{r}$

MEANING ATTENUATION WITH r
AS $r \rightarrow \infty$ PLANE WAVE!
AS $r \rightarrow 0$ BLOWS UP

$$P(r, t) = \frac{1}{r} f\left(t - \frac{r}{c}\right) + \frac{1}{r} g\left(t + \frac{r}{c}\right)$$

f & g ARE CONTINUOUS BUT AT
LEAST A SECOND DERIVATIVE

MEANING

$t - \frac{r}{c} \rightarrow$ MOVES AWAY

$t + \frac{r}{c} \rightarrow$ MOVES INWARD

CAUSAL - HIT IT BEFORE IT RESPONDS

$$\theta \approx 0$$

$\frac{\Delta r}{c}$ TIME FROM DISTURBANCE TO
LISTENER

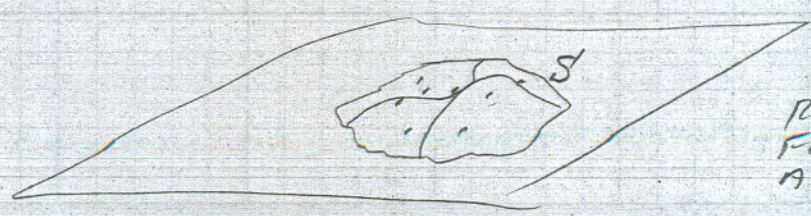
$\frac{1}{r}$ ATTENUATION WITH TIME

MORE ADVANCED CONCEPT - FUNDAMENTAL
SOLUTION

$$\iint_S P(r, t) ds \text{ TO USE}$$

SOLUTION FROM POINT SOURCE

OR WITH AMPLITUDE

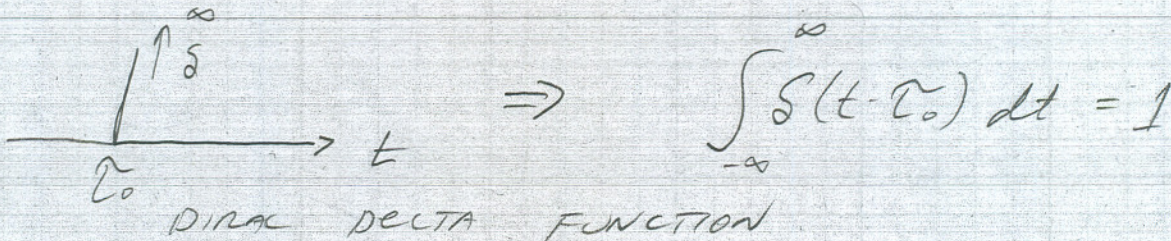


RADIATED
FIELD FROM
A SOURCE

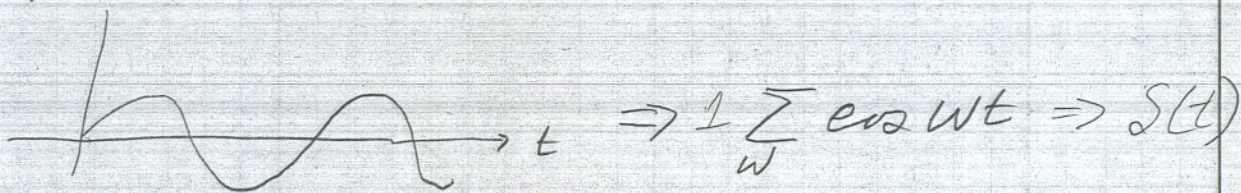
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We will focus on harmonic waves -
Will return but 2 fundamental solutions -



DIRAC DELTA FUNCTION



HARMONIC SOLUTION

ADD UP ALL HARMONIC SOLUTIONS
GET A $\delta(t)$ AT 0

FOR OUR SPHERICAL CASE

$$\frac{\partial^2 p}{\partial r^2} + \frac{2}{r} \frac{\partial p}{\partial r} = \frac{1}{c^2} \frac{\partial^2 p}{\partial t^2}$$

$$\text{OR } \frac{\partial^2(r p)}{\partial r^2} = \frac{1}{c^2} \frac{\partial^2(r p)}{\partial t^2}$$

$$\tilde{p} = \frac{\tilde{A}}{r} e^{i(\omega t - kr)}$$

RECALL Euler's IDENTITY

$$e^{i\theta} = \cos \theta + i \sin \theta$$

AT A LARGE DISTANCE r



$r \rightarrow \infty$
PLANE WAVE SOLUTION

Thus we see that a spherical spreading wave becomes the plane wave as $r \rightarrow \infty$

Using FEXT NOTATION NOTE

u IS VELOCITY NOT DISPLACEMENT

WE REITERATE

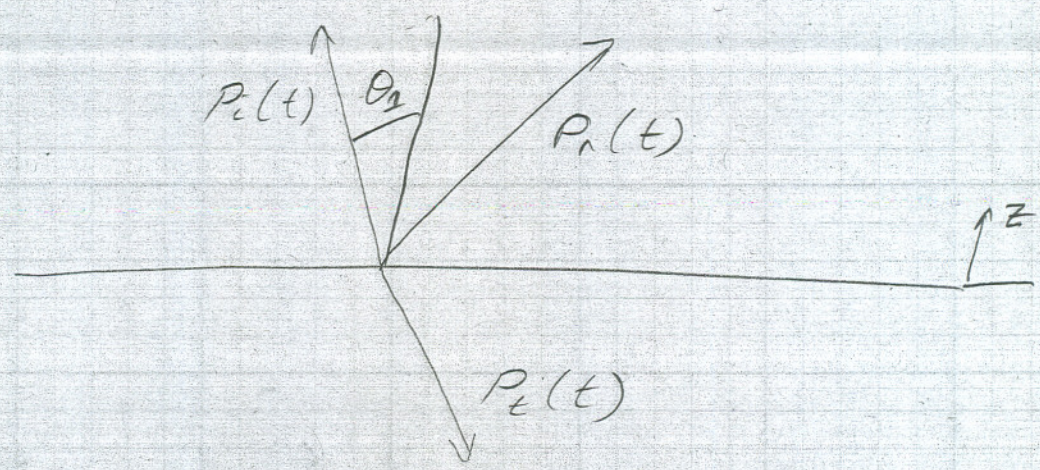
$$\frac{\partial u}{\partial t} = -c \frac{\partial u}{\partial x}$$

$$\frac{\partial p}{\partial x} = \rho \frac{\partial u}{\partial t} \quad F = ma$$

$$\frac{\partial p}{\partial x} = \boxed{\rho c} \frac{\partial u}{\partial t}$$

ACOUSTIC IMPEDANCE -

ACOUSTIC IMPEDANCE IS THE CHARACTERISTIC ABILITY TO TRANSFER ENERGY BETWEEN MEDIA -



B.C. pressure is continuous
NORMAL COMPONENTS OF VELOCITY ARE EQUAL -

$$p_i(t) + p_r(t) = p_t(t)$$

$$u_{z_i}(t) + u_{z_r}(t) = u_{z_t}(t)$$

$$p = \rho c u$$

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$$\frac{\sin \theta_2}{c_2} = \frac{\sin \theta_1}{c_1} \quad \text{Snell's Law}$$

THE VERTICAL (z) VELOCITIES

$$\frac{v_{c_1 c_1}}{c_2 c_2}$$

$$u_{zi}(t) = u_i(t) \cos \theta_1$$

$$u_{zn}(t) = u_n(t) \cos \theta_1$$

$$u_{zt}(t) = u_t(t) \cos \theta_2$$

Pressure $p = -(\rho c)u$

THEN

$$u_{zi}(t) = \frac{p_i(t)}{\rho_i c_i} \cos \theta_1$$

$$u_{zn}(t) = \frac{p_n(t)}{\rho_i c_i} \cos \theta_1$$

$$\theta_i = \theta_r$$

$$u_{zt}(t) = \frac{p_t(t)}{\rho_t c_t} \cos \theta_2$$

THEN

define $\frac{p_r}{p_i} = R_{12}$ REFLECTION

$\frac{p_t}{p_i} = T_{12}$ TRANSMISSION

$$p_t = p_i(t - z/c_2)$$

$$p_i = p_t(t - z/c_1)$$

$$p_r = p_n(t + z/c_1)$$

LOOK

AT VELOCITY

$$u_{zi} + u_{zn} = u_{zt}$$

$$u_i = p_i \cos \theta_1 (kx - \omega t) \text{ or } p_i f(kx - \omega t)$$

$$u_r = R_{12} p_i f(kx - \omega t)$$

$$u_t = T_{12} p_i f(kx - \omega t)$$

13,782 500 SHEETS FULLER 8 SQUARE
42,381 50 SHEETS FULLER 8 SQUARE
42,382 100 SHEETS FULLER 8 SQUARE
42,383 100 SHEETS FULLER 8 SQUARE
42,389 100 SHEETS FULLER 8 SQUARE
42,390 200 SHEETS FULLER 8 SQUARE
42,395 200 RECYCLED WHITE 8 SQUARE
Made in U.S.A.



OR $1 + R_{12} = T_{12}$ PRESSURE EQU

VELOCITY $u_2 + u_n = u_t$

$$(\rho_2 c_2 - \rho_1 c_2 R_{12}) \cos \theta_1 = \rho_1 c_1 T_{12} \cos \theta_1$$

SINCE

$$\frac{\rho_2}{\rho_1 c_1} \cos \theta_1 + R_{12} \left(\frac{-\rho_2}{\rho_1 c_1} \right) \cos \theta_1 = T_{12} \frac{\rho_2}{\rho_1 c_2} \cos \theta_1$$

← MOVING -Z

$$R_{12} = \frac{\rho_2 c_2 \cos \theta_1 - \rho_1 c_1 \cos \theta_2}{\rho_2 c_2 \cos \theta_1 + \rho_1 c_1 \cos \theta_2}$$

$$T_{12} = \frac{2 \rho_2 c_2 \cos \theta_1}{\rho_2 c_2 \cos \theta_1 + \rho_1 c_1 \cos \theta_2}$$

$$\theta_2 = \sin^{-1} \left(\frac{c_2}{c_1} \sin \theta_1 \right)$$