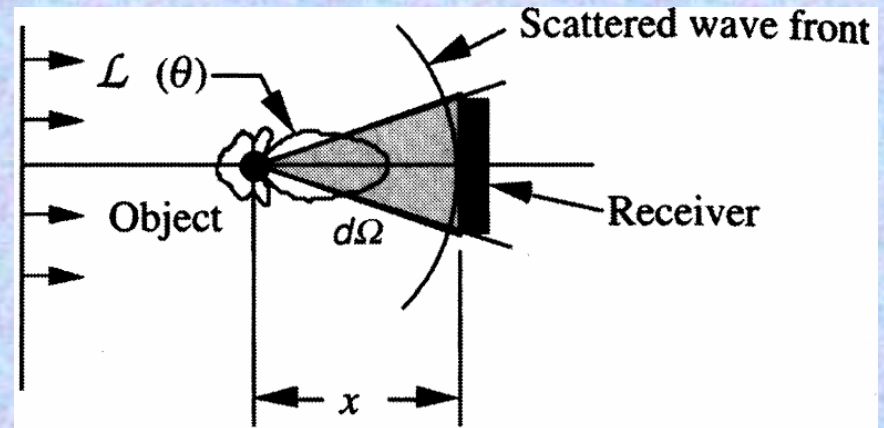


## Scattering of sound - Part II- scattering from a sphere

The problem:

- An incident plane-parallel wave
- Possible absorption
- A scattered wave



Medwin et al., 2005

Scattering: combined action of reflection, refraction and diffraction from objects with differing sound speed.

References:

Anderson, V. C., 1950: Sound scattering from a fluid sphere, *JASA*, 22, 426-431.

Faran, J. J., 1951: Sound scattering by solid cylinders and spheres, *JASA*, 23, 405-418.

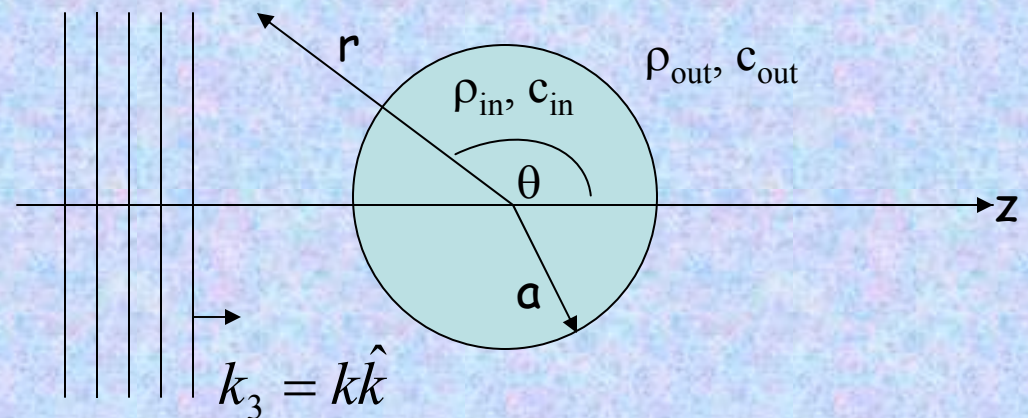
Medwin et al., 2005, Ch. 6.

Williams, E. G., 1999: *Fourier Acoustics*, Academic Press, 306pp.

## Formulation of the problem:

Plane-wave insonification (neglect attenuation within the medium):

$$p_i(\vec{r}, t) = A_0 \exp\{i(\vec{k} \cdot \vec{x} - \omega t)\} = A_0 \exp\{i(kr \cos \theta - \omega t)\}$$



The problem is azimuthally symmetric, which facilitate its formulation in spherical coordinates.

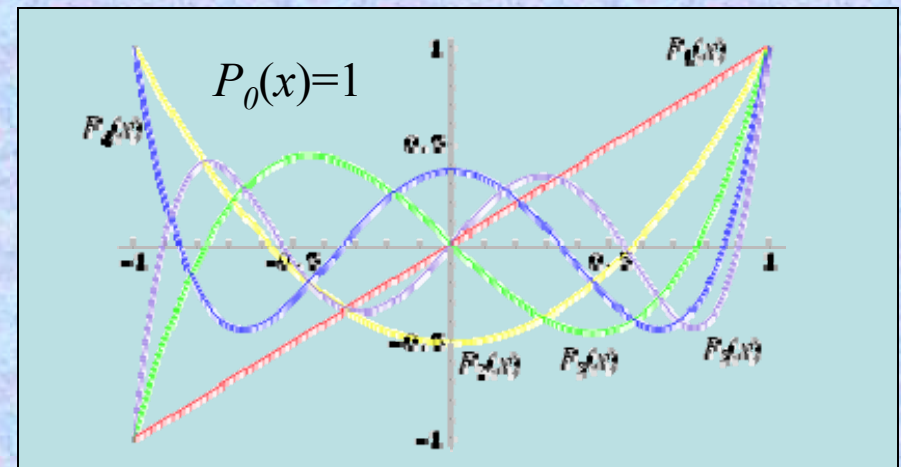
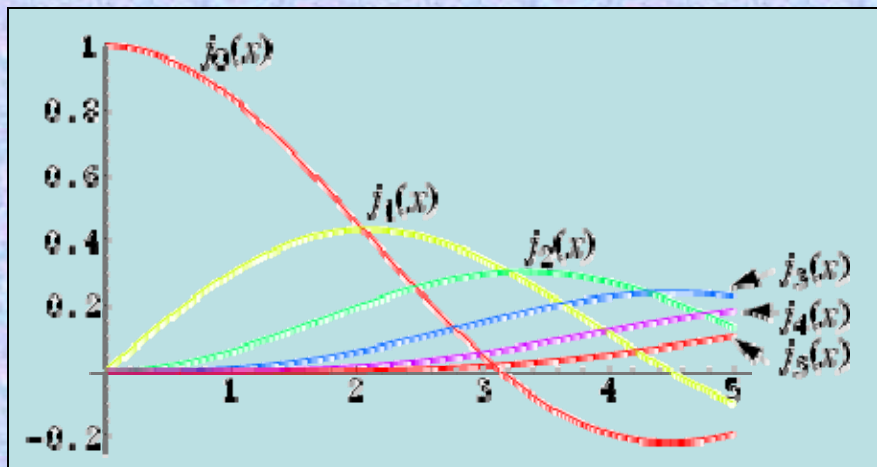
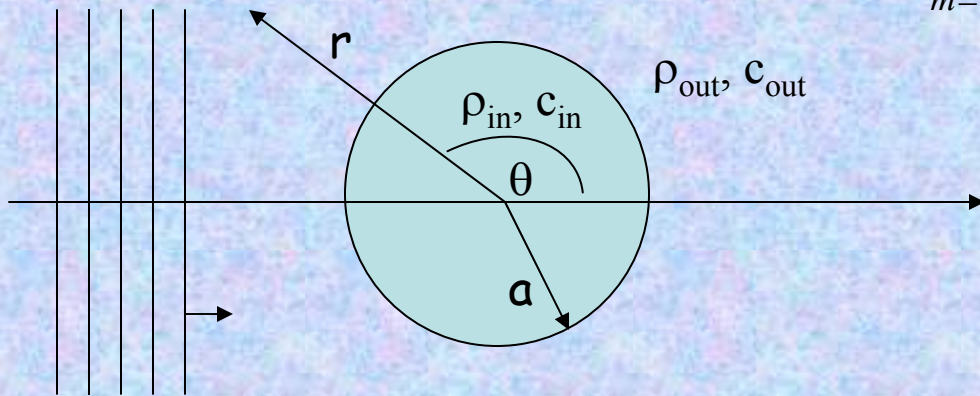
The scattered wave also obeys:

$$\nabla^2 p_s = \frac{1}{c^2} \frac{\partial^2 p_s}{\partial t^2}$$

# Formulation of the problem (scattering by a fluid sphere):

We need to express the incident field in spherical coordinates which for azimuthal symmetry is given by:

$$p_i(\vec{x}, \theta, \omega) = A_0 \exp(ikr \cos \theta - \omega t) = A_0 \sum_{m=0}^{\infty} (-i)^m (2m+1) j_m(kr) P_m(\cos \theta) e^{-i\omega t}$$



BCs: radial velocity and pressure are continuous across the boundary.

$$u_i(a^+, \theta) + u_s(a^+, \theta) = u_{in}(a^-, \theta)$$

$$p_i(a^+, \theta) + p_s(a^+, \theta) = p_{in}(a^-, \theta)$$

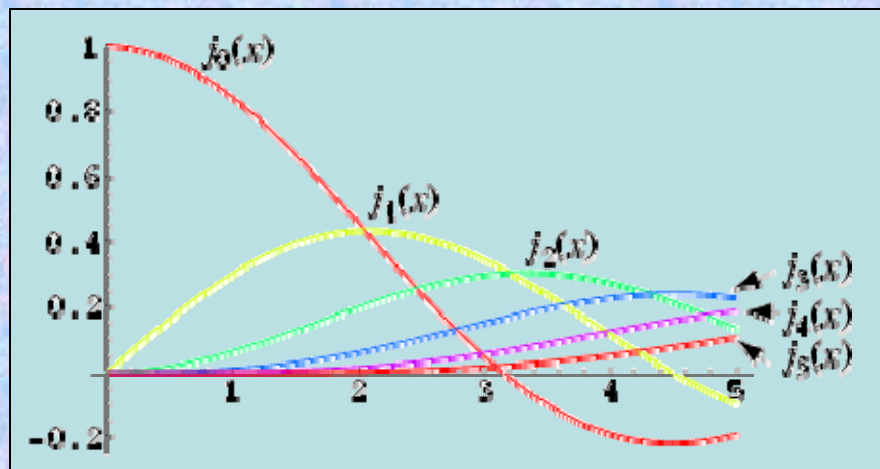
$$\omega_{in} = c_{in} k_{in} / 2\pi = c_{out} k / 2\pi = \omega$$

The scattered sound out of the sphere and the pressure within the sphere are given as the general axisymmetric solution for the wave equation that vanishes at  $r \rightarrow \infty$ ):

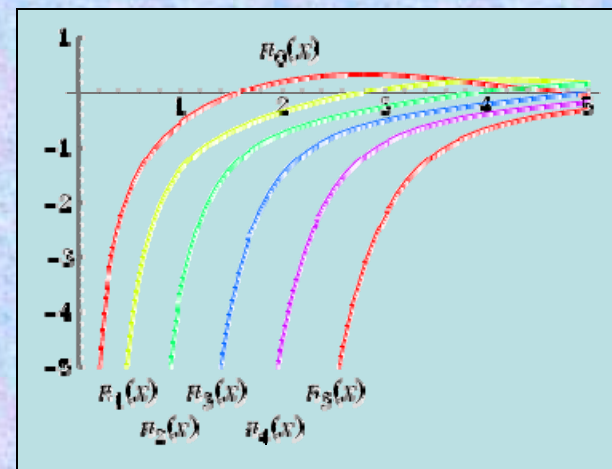
$$p_s(r, \theta) = \sum_{m=0}^{\infty} \boxed{C_m} h_m^{(1)}(kr) P_m(\cos \theta), \quad h_m^{(1)}(x) = j_m(x) + in_m(x)$$

$$p_{in}(r, \theta) = \sum_{m=0}^{\infty} \boxed{B_m} j_m(k_{in} r) P_m(\cos \theta)$$

$\boxed{\phantom{C_m}}$  ← Unknowns at this stage



+i



Euler equation:

$$\rho \frac{\partial u_r}{\partial t} = -\frac{\partial p}{\partial r} \Rightarrow u_r = -\frac{i}{\rho c} \frac{\partial p}{\partial (kr)}$$

→ 2 more equations (in and out) for a total of 4 to solve for complex  $B_m$  and  $C_m$ .

$$C_m = \frac{-A_0 (-i)^m (2m+1)}{(1+iD_m)}, \text{ where:}$$

$$D_m \equiv \frac{\{mj_{m-1}(k'a) - (m+1)j_{m+1}(k'a)\}n_m(ka) - \{j_m(k'a)n_{m-1}(ka) - (m+1)n_{m+1}(ka)\}gh}{\{mj_{m-1}(k'a) - (m+1)j_{m+1}(k'a)\}j_m(ka) - \{mj_{m-1}(ka) - (m+1)j_{m+1}(ka)\}gh}$$

and  $g = \rho_{\text{in}}/\rho_{\text{out}}$ ,  $h = c_{\text{in}}/c_{\text{out}}$ .

Thus the scattered wave at any point outside the sphere is given by:

$$p_s(r, \theta) = -A_0 \sum_{m=0}^{\infty} (-i)^m \frac{(2m+1)}{(1+iD_m)} h_m^{(1)}(kr) P_m(\cos \theta) e^{-i\omega t}$$

Far field:  $r \rightarrow \infty$ :

$$j_m(kr) \xrightarrow{kr \rightarrow \infty} \frac{1}{kr} \cos\left(kr - (m+1)\frac{\pi}{2}\right)$$

$$n_m(kr) \xrightarrow{kr \rightarrow \infty} \frac{1}{kr} \sin\left(kr - (m+1)\frac{\pi}{2}\right)$$

$$\Rightarrow (-i)^m h_m^{(1)} \xrightarrow{kr \rightarrow \infty} (-1)^m \frac{1}{kr} e^{ikr}$$

→ Far field solution for scattered wave:

$$p_s(r, \theta) = A_0 \frac{ie^{-i(kr-\omega t)}}{kr} \sum_{m=0}^{\infty} (-1)^m \frac{(2m+1)}{(1+iD_m)} P_m(\cos \theta)$$

Normalizing factor- let assume the whole intensity impinging on the sphere is radiated out uniformly:

$$A_0^2 \pi a^2 = |p_{s,norm}|^2 4\pi r^2 \rightarrow |p_{s,norm}| = \frac{A_0 a}{2r}$$

Define 'reflectivity':

$$R(\theta) \equiv \frac{|p_s(r, \theta)|}{|p_{s, norm}(r, \theta)|} = \frac{2}{ka} \left| \sum_{m=0}^{\infty} (-1)^m \frac{(2m+1)}{(1+iD_m)} P_m(\cos \theta) \right|$$

and for backscattering (sometimes called 'form function',  $f_{\infty}$ ):

$$R = R(\theta = 0) = \frac{2}{ka} \left| \sum_{m=0}^{\infty} (-1)^m \frac{(2m+1)}{(1+iD_m)} \right|$$

with respect to 'differential scattering cross-section' and 'scattering length':

$$\Delta\sigma_s(\theta, k) \equiv |L(\theta, k)|^2 \equiv \frac{|p_s|^2 r^2}{|p_i|^2} = \frac{R(\theta)^2 a^2}{4}$$

Normalized power (the amount of power the scatterer diverts):

$$\Pi_{norm} = \frac{\int |p_s(r, \theta)|^2 ds}{A_0^2 \pi a^2} = \dots = \frac{4}{(ka)^2} \left| \sum_{m=0}^{\infty} \frac{(2m+1)}{(1+D_m)} \right| = \frac{\sigma_s}{\pi a^2}$$

# Scattering from a sphere:

Limiting cases: I. Rayleigh scattering ( $ka \ll 1$ )

$$j_m(x) \xrightarrow{x \rightarrow 0} x^m / 1 \cdot 3 \cdots (2m+1)$$

$$n_m(x) \xrightarrow{x \rightarrow 0} -1 \cdot 3 \cdots (2m-1) / x^{m+1}$$

$$R(\theta) \rightarrow 2(ka)^2 \left| \frac{1-gh^2}{3gh^2} + \cos \theta \frac{1-g}{1+2g} \right|$$

$$gh^2 = E_{in}/E_{out}$$

$$g = \rho_{in}/\rho_{out}, \quad h = c_{in}/c_{out}$$

$$\Rightarrow \Delta\sigma_s, \sigma_s, \Pi_s \propto (ka)^4 \quad \text{Very sensitive to frequency!}$$

II. Geometric acoustics ( $ka \gg 1$ ):

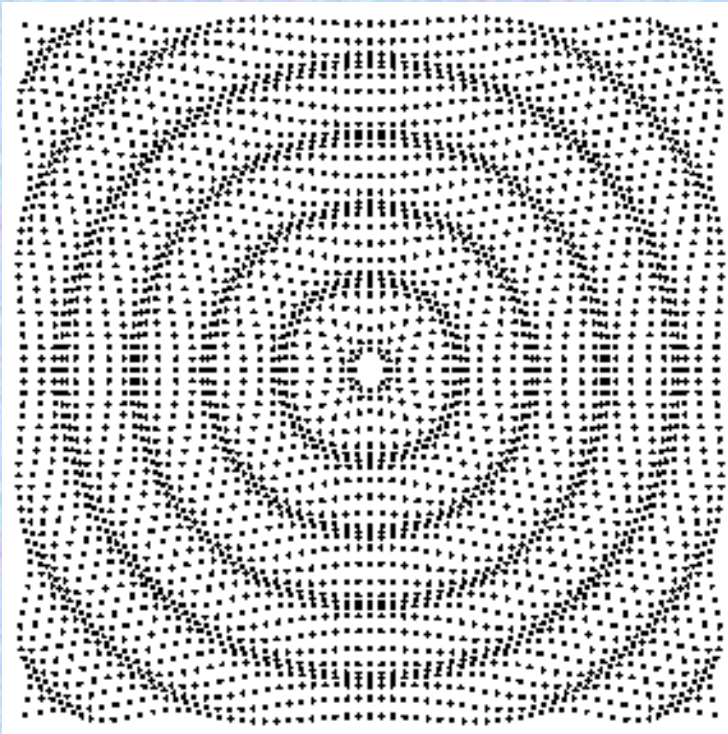
$$\Pi_{norm} = \frac{\int |p_s(r, \theta)|^2 ds}{A_0^2 \pi a^2} = \dots = 2 = \frac{\sigma_s}{\pi a^2}$$

Independent of  $ka$  !!!

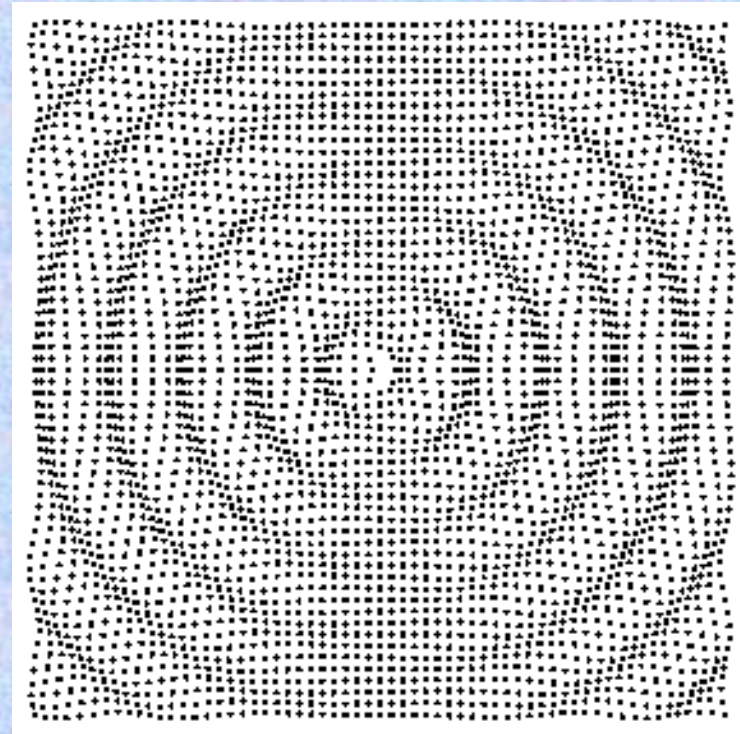
???



Monopole:



Dipole:



Animations courtesy of Dr. Dan Russell, Kettering University

# Scattering from a sphere:

Three types of particles:

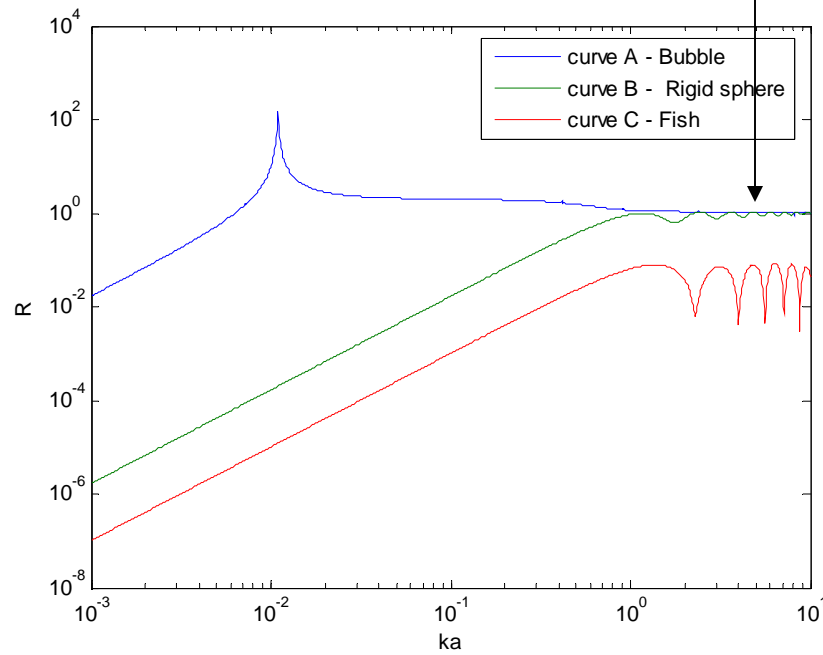
I. Rigid sphere,  $c_{in} \gg c_{out}$ ,  $\rho_{in} \gg \rho_{out}$ . ← sediments, calibration targets.

II. Generic 'Marine particle (fish, zooplankton)':  $1.08 > g > 1.02$ ,  $1.033 > h > 1.00$

III. Bubble ( $\rho_{in}/\rho_{out} \sim 1/1000$ ,  $c_{in}/c_{out} \sim 1/5$ )

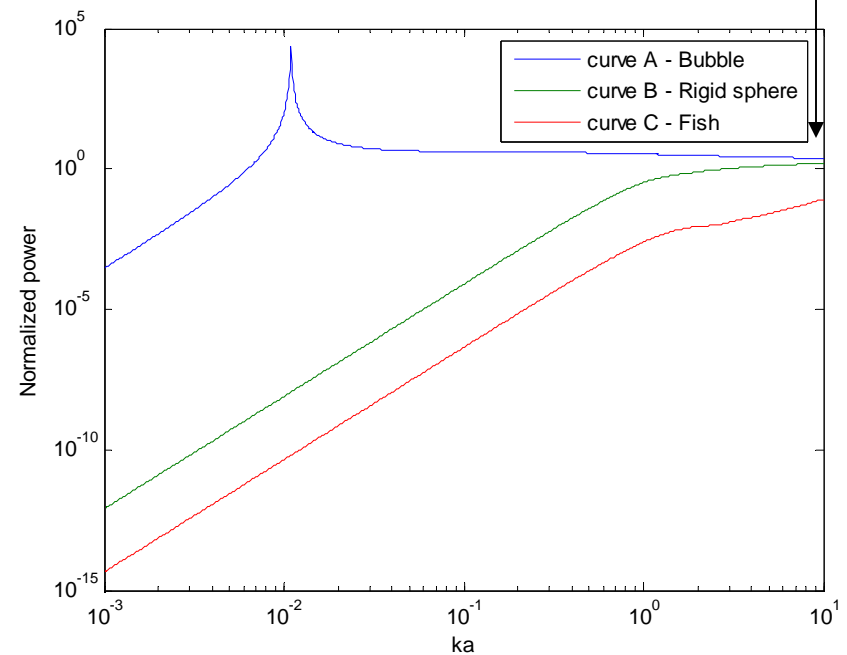
Form-function (reflectivity)

$R=1$



Normalized power ( $\Pi_{norm}$ )

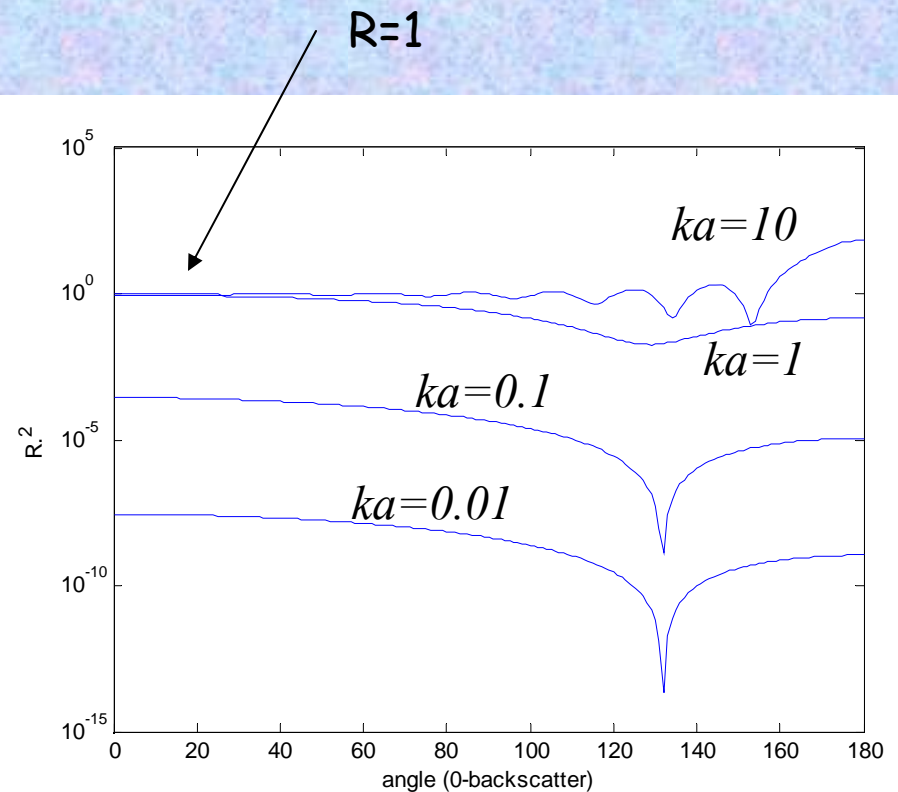
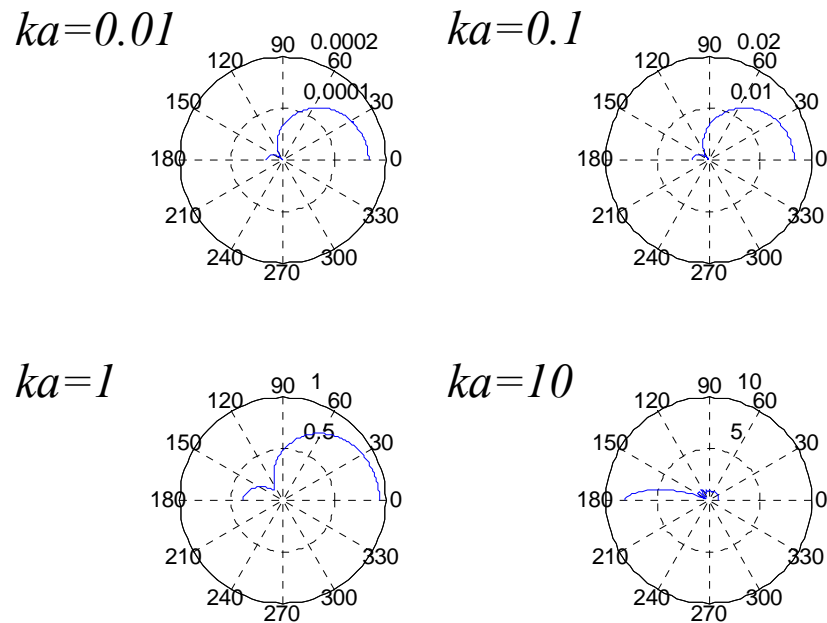
$\Pi_{norm} = 2$



# Scattering from a sphere:

I. Rigid sphere,  $c_{in} \gg c_{out}$ ,  $\rho_{in} \gg \rho_{out}$ .

directional distribution:



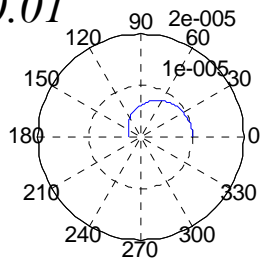
Note: our treatment has neglected shear waves which may exist (Faran, 1951).

# Scattering from a sphere:

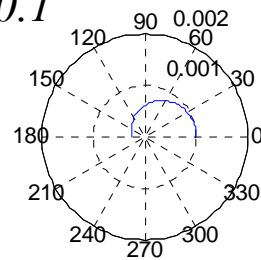
II. Generic 'Marine particle (fish, zooplankton)':  $1.08 > g > 1.02$ ,  $1.033 > h > 1.007$

directional distribution:  
Fish

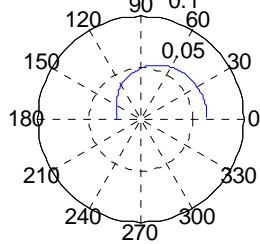
$ka=0.01$



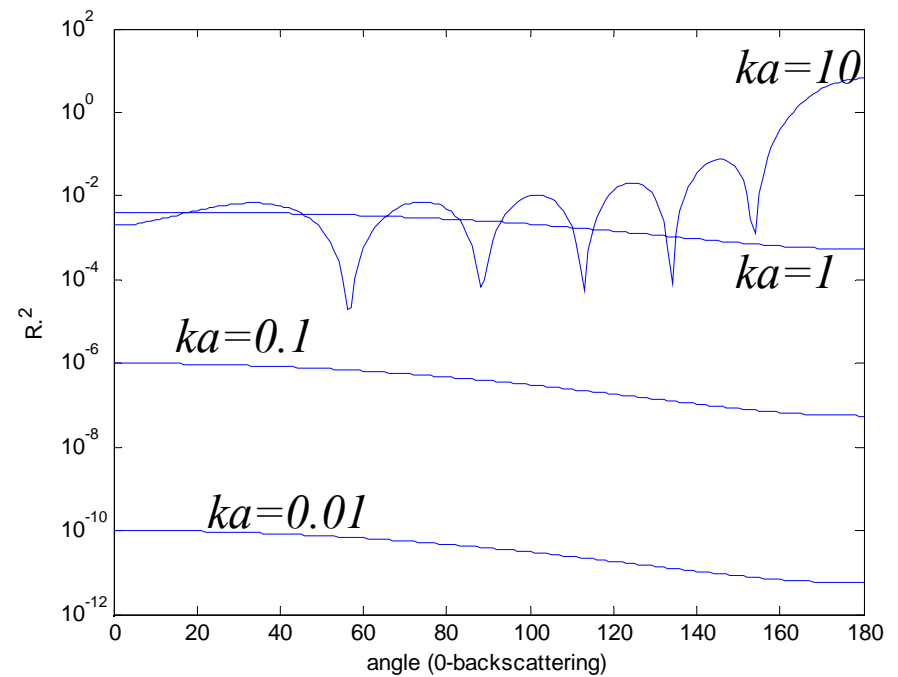
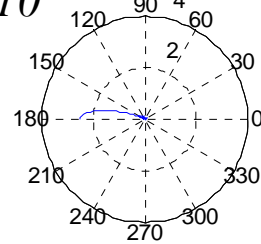
$ka=0.1$



$ka=1$



$ka=10$

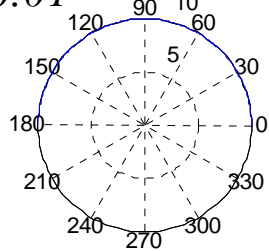


# Scattering from a sphere:

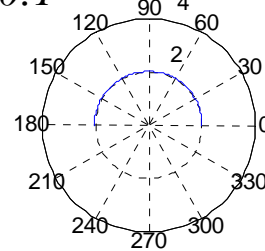
III. Bubble ( $\rho_{in}/\rho_{out} \sim 1/1000$ ,  $c_{in}/c_{out} \sim 1/5$ )

directional distribution:

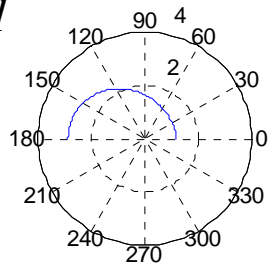
$ka=0.01$



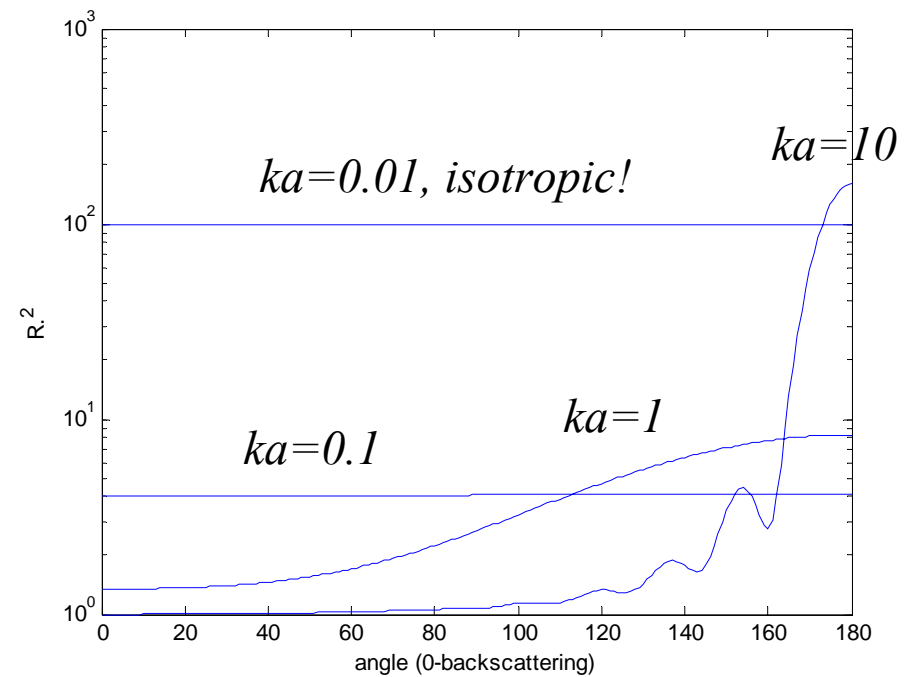
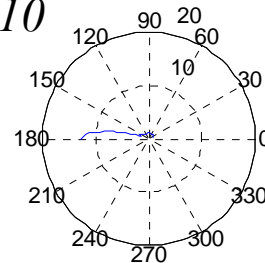
$ka=0.1$



$ka=1$



$ka=10$



How do we get an absorption, scattering and attenuation coefficient from all this?

$$\sigma_a \equiv \frac{\Pi_a}{I_i}, \sigma_s \equiv \frac{\Pi_s}{I_i}, \sigma_e = \sigma_a + \sigma_s$$

$$a(k) = \sum_i \sigma_a(k, a_i) N_i$$

$$b(k) = \sum_i \sigma_s(k, a_i) N_i$$

$$\gamma(k) = \sum_i \sigma_e(k, a_i) N_i$$

Beer-Lambert-Bouguer's law: in a dilute system with a variety of different substances, the absorption, scattering and attenuation of the different substances sum together to provide the absorption, scattering and attenuation of the medium:

$$I(r) = I(r=0) \exp\left(-r \sum_j \gamma_j\right)$$