

### 4.3 Koschmieder's Equation for Radiance

A classical problem of radiative transfer theory in either the atmosphere or in the sea is to determine the apparent radiance of an object as seen along a path of sight,  $O(x, 0)$  which lies in a homogeneous and uniformly lighted region of an optical medium.

Specifically, the problem is to determine the apparent radiance  $N_r(z, \sim)$  given  $a$  and  $a_0$  along  $O(x, \sim)$ , and  $N_0(x, \sim)$  at the initial end point  $x$  of the path  $\sim$  along with the fact that each point of  $O(x, \sim)$  is irradiated by the same radiance distribution (which may, however, depend arbitrarily on  $\sim$ ). This situation (or some reasonable approximation of it) arises often in the atmosphere and the sea, notably along horizontal paths of sight,

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and the reader should be able to cite many personally observed instances of it. Koschmieder studied this classical setting in detail, and in 1924 published in [141] his analytic expression for  $N_r(z, 0)$  which was derived after lengthy preliminaries and under the radiometric conditions stipulated above. We turn now to a modern derivation of the expression for  $N_r(z, \sim)$ .

Returning to (1) of Sec. 3.15 we assume  $a$  and  $a_0$  are independent of  $z$  along  $O(x, \sim)$ . Then

$$N_r(z, \sim) = e^{-az} N_0(x, \sim)$$

where "a" denotes the assumed fixed value of the volume attenuation function along  $O(x, \sim)$ . Furthermore, since the radiance distribution  $N_r(z, \sim)$  is independent of  $z$  along the path then  $N_r(z, E)$

is also independent of  $z$  along the path and we shall abbreviate this fixed value by " $N^*$ ". Equation (1) of Sec. 3.15 then reduces to:  

$$N_r(z, E) = a N_0(x, E) e^{-az} + N^*$$

$$e^{-az} \frac{dN_r}{dz} + N_r = a N_0 + N^*$$

and with the abbreviations for  $N_r(z, \sim)$  and " $N_0$ " for  $N_0(x, \sim)$ , this simplifies immediately to

where we have written:

$$f_1$$

for  $N^*/a$

Equation (1) is Koschmieder's equation which relates apparent radiance  $N_r$  to  $N_0$  on a path  $O$  in an optical medium along which  $a$  and  $a_0$  are constant values and along which the value  $N^*$  of the path function is constant. The radiance  $N$  is called the equilibrium radiance for  $O_r$ . The significance &

of  $N$  is seen by letting  $r \rightarrow \infty$  in (1), or alternately by contemplating the integrodifferential equation for  $N_r$  associated with  $O_r$  as given in (3) of Sec. 3.15

$$dN_r$$

Under our present assumptions, (3) is a relatively innocuous first order differential equation in which  $a$  and  $N^*$  are constants and  $N_r$  is the unknown function. Using (2) we can rewrite (3) as

SEC. 4 . 3 KOSCHMIEDER' S EQUATION  $\frac{dN_r}{dr} = \frac{N_r(N^* - N_r)}{r}$  from which we can immediately read the physical significance of  $N_q$ : If  $N_r < N_q$  at a point on the path, then  $dN_r/dr > 0$ , i. e.,  $N_r$  is increasing at that point. In general,  $N$  always tends toward the fixed radiance  $N_q$ , and  $dN_r/dr = 0$  if and only if  $N_r = N_q$ . Therefore  $N_q$  takes on the aspect of an equilibrium value (in an ever' day sense) toward which the values  $N_r$  unceasingly tend. The equilibrium radiance  $N_q$  is often observable over long horizontal uniformly lighted paths through a homogeneous natural aerosol or hydrosol.

It should be observed that the derivation of (1) places no conditions on the orientation or the location of the path ( $r$  in an optical, medium. The essential point to observe in the derivation is that (1) follows from (1) of Sec. 3.15 upon assuming only that  $a$ ,  $a$  and  $N^*$  are constant long  $r$ . This leaves  $r$  free to be vertical, inclined, or horizontal, as the case may be. An interesting example of (1) for inclined paths of sight in the atmosphere may be obtained from the results in [71].